



## Surface-Energy Generator of Single-Walled Carbon Nanotubes and Usage in a Self-Powered System

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Surface energy plays an important role in surface physics,<sup>[1,2]</sup> biophysics,<sup>[3,4]</sup> surface chemistry,<sup>[5,6]</sup> and catalysis.<sup>[7]</sup> A gradient of surface energy between a solid and liquid interface can induce transport of liquids<sup>[8-11]</sup> and water running uphill,<sup>[12]</sup> which is important for DNA analysis devices.<sup>[13]</sup> Due to the 2D nature and relatively few molecules or atoms involved, the density of surface energy is usually quite small, which is impractical for utilizing surface energy as an energy source. Nevertheless it is attractive to use surface energy at the nanoscale because of the lower power consumption for nanodevices and the higher specific surface area for nanomaterials.<sup>[14–18]</sup> In this work, we demonstrate an effective design of single-walled carbon nanotubes (SWNTs) to harvest surface energy of ethanol and convert it into electricity. In this ethanol-burner-like design, an open-circuit voltage ( $V_{oc}$ ) can be obtained as a result of ethanol flow in the capillary channels formed among SWNTs driven by surface tension. The Voc remains constant as long as there is ethanol from the source. The maximum power can be up to  $\sim$ 1770 pW per device and can serve as a self-powered system to drive a thermistor. Meanwhile, the performance (the inducing rate of  $V_{\rm oc}$ , the value of  $V_{\rm oc}$ , and the output power) can be significantly enhanced by the Marangoni effect.[19]

SWNTs were synthesized by floating catalytic chemical vapor deposition and treated by diamond wire drawing dies, <sup>[20–22]</sup> which results in well-aligned individual SWNTs (see Supporting Information S1). The resulting SWNT rope (~25.0 mm (length) × 0.6 mm (diameter), Fig. 1a) is connected to electrodes of aluminum film, forming a suspended structure on a glass

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slide. The device is measured by a Keithley 4200-SCS, semiconductor characterization system, (voltage resolution  $1 \mu V$ ) and the dynamic characteristics of the open-circuit voltage ( $V_{\rm oc}$ ) are monitored while adding ethanol (MOS grade, 99.9%) to the beaker (Fig. 1b, see Supporting Information S2).

In an open beaker, no obvious  $V_{oc}$  is observed at the beginning (angle 40°, Fig. 2a). When the ethanol level reaches the SWNT rope, the  $V_{oc}$  begins to increase. The increase of  $V_{oc}$  is almost linear from zero to 200 µV for the first 240 s, then the  $V_{oc}$ saturates gradually at 219 µV where is remains constant over 6 h as shown in Figure 2a.  $V_{oc}$  can remain constant as long as the ethanol level is contacting the SWNT rope. When the beaker is covered as indicated by region 2 (Fig. 2b),  $V_{oc}$  will gradually decrease back to the original value. This process can be repeated



**Figure 1.** The SWNT device and schematic layout of the experimental setup. a) An image of the device with a suspended SWNT rope. After treating with diamond wire drawing dies, the SWNT rope has a diameter of ~0.6 mm and length of ~25.0 mm and is connected to the electrodes. b) When measuring, the device is placed into a beaker with an angle between the SWNT rope and the ethanol level. Dynamic characteristics of  $V_{\rm oc}$  are monitored when ethanol is added into the beaker.





**Figure 2.** Dynamic characteristics of  $V_{oc}$  at  $\alpha = 40^{\circ}$  of the SWNT rope. a) When ethanol is added into the beaker, no  $V_{\rm oc}$  is observed until the ethanol level reaches the SWNT rope. The  $V_{\rm oc}$ increases to a saturated value of  $219 \,\mu$ V in about 300 s and remains constant for about 6 h. b) In an open beaker, the  $V_{oc}$  will begin to increase when the ethanol level reaches the SWNT rope and tend to a saturated value (region 1). The  $V_{oc}$  will decrease gradually back to the original value when the beaker is covered as indicated by region 2. These characteristics can be reproduced. c) When reaching the saturated value of  $V_{oc}$ , the  $V_{oc}$  does not change even if more ethanol is added (arrow 1). After certain amount of ethanol is added, the  $V_{oc}$  begins to decrease. At both points 2 and 3, 3.0 mL ethanol is added. At points 4, 5, and 6, 9.0 mL ethanol is added. At point 7, the SWNT rope is almost immersed in ethanol and no obvious change of  $V_{oc}$  is found. A similar result is obtained at point 8 when SWNT is totally immersed. According to the diameter of the beaker (68 mm) and the angle between the SWNT rope and the ethanol level (40°), the L of the SWNT rope is 14.1 mm. d) Dependence of  $V_{oc}$  and L on  $\alpha$  of the SWNT rope. Both  $V_{oc}$  and L are closely related to  $\alpha$  of the SWNT rope. The angles are 90°, 60°, 40°, and 17° for V<sub>oc</sub> measurements. The maximum power can be up to  $\sim$ 1770 pW per device (17°). L is 14.1, 10.6, and 7.6 mm for angles of  $90^{\circ}$ ,  $60^{\circ}$ ,  $40^{\circ}$ , respectively.

and the characteristics of varying  $V_{oc}$  can be well reproduced (Fig. 2b).

When reaching the saturated value,  $V_{oc}$  does not change until a certain amount of ethanol is added into the beaker. When 3.0 mL of ethanol is added into the beaker as indicated by the time points 2 and 3 in Figure 2c, a drop of  $\mathit{V}_{oc}$  (~20  $\mu V)$  is found. The addition of 9.0 mL of ethanol is carried out at points 4, 5, and 6. At this time, the ethanol level almost reaches the top electrode. The addition of ethanol does not cause an obvious change for time points 7 and 8, as shown in Figure 2c. These results suggest a characteristic length (L) for the device. When the length of the SWNT rope over the ethanol level is longer than L, the induced  $V_{oc}$ always takes the maximum value and the addition of ethanol into the beaker does not cause any change of V<sub>oc</sub>. If the length of the SWNT rope over the ethanol level is equal to L, the induced  $V_{oc}$ begins to decrease with additional ethanol. The L can be calculated using  $V/(\pi r^2 \sin \alpha)$ , where V is the total volume of ethanol that causes changes in  $V_{oc}$ , r is the radius of the beaker, and  $\alpha$  is the slant angle of the SWNT. The maximum  $V_{\rm oc}$  and L are closely related to  $\alpha$  and device-dependent. For the device reported, the relationships between  $V_{\rm oc}$ , *L* and  $\alpha$  are plotted in Figure 2d.



More interesting phenomena can be found if the SWNT rope is pretreated. Before measurements for  $V_{\rm oc}$ , the device can be fixed where droplets of deionized water (Millipore,  $18 \,\mathrm{M}\Omega \,\mathrm{cm}^{-1}$ ) fall on the SWNT rope. After this treatment for several minutes, the device is transferred into the beaker for further studies (Fig. 1b). When ethanol is added into the beaker and its level reaches the SWNT, a very sharp increasing of  $V_{\rm oc}$  is observed (Fig. 3b). The  $V_{\rm oc}$  jumps to 853  $\mu$ V in less than 1 s (inset of Fig. 3b), representing a significant increasing rate compared to that in Figure 3a. The  $V_{oc}$ remains at this value for about 50 s and decreases gradually to a value of  $\sim 223 \,\mu V$ , which is quite close to that when using pure ethanol (Fig. 3a).

What is the mechanism of these interesting observations? A possible electrochemical potential difference at the metal/SWNT interface for  $V_{oc}$  can be excluded for two reasons. Firstly, when ethanol level reaches the SWNT rope, a linear increasing behavior of  $V_{oc}$  is observed instead of a sharp one, which is expected if there exists an electrochemical potential difference for the faster response time of hundreds of milliseconds to several seconds (Fig. 2a).<sup>[23]</sup> Secondly, if the beaker is covered, the  $V_{oc}$  decreases to a value of nearly zero, which indicates that the electrochemical potential difference at the metal/SWNT interface is negligible.

When the ethanol level reaches the SWNT rope, surface tension pulls the ethanol molecules up along the channels formed among individual SWNTs, which have dimensions around hundreds of nanometers, because of the ethanol contact angle of  $\sim 0^{\circ}$  (Supporting

Information, Fig. S1b and S3). The moving ethanol molecules along the SWNT can induce a  $V_{\rm oc}$  because of the coupling between the charge carriers of the SWNTs and the flowing molecules at the interface.<sup>[22,24,25]</sup> At capillary equilibrium, however, ethanol no longer moves and hence no  $V_{\rm oc}$  should be observed. This is in contradiction with our experimental results. Based on the device design and experimental results, we propose an ethanol-burnerlike model as shown in Figure 3c and d. When ethanol molecules climb up along the channels among the SWNTs, these molecules also evaporate from the rising liquid at the same time. The average length that ethanol molecules move along the channel before they evaporate is L (Fig. 3c). When certain ethanol molecules pass L and evaporate at the top, the same quantity of ethanol will climb up along the channels among the SWNTs from the bottom. This forms a steady, dynamic, and directional ethanol flow and results in the constant Voc that can be obtained as long as the SWNT rope is immersed in ethanol (Fig. 2a). L cannot be visually observed in our experiment because the SWNT rope is black whether it is wetted by ethanol or not.

Poiseuille's law can be used to estimate linear velocity of this novel ethanol laminar flow at the steady state.<sup>[26,27]</sup> The linear



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**Figure 3.** Comparison between the SEG and the Marangoni-enhanced SEG and diagrams showing the mechanisms. a) When the SWNT rope is immersed by ethanol, the  $V_{oc}$  increases and reaches a value of ~219  $\mu$ V in about 300 s. b) Once the SWNT is pretreated with water, the  $V_{oc}$  jumps to a value of 853  $\mu$ V in less than 1 s, indicating a much higher inducing rate and larger value of  $V_{oc}$  of the SEG by the Marangoni effect. Usually the  $V_{oc}$  maintains at the maximum value for tens of seconds and decreases gradually to a value similar to that of the pure ethanol SEG. c) Once the ethanol level reaches the SWNT rope, the ethanol climbs along the channels among the SWNTs due to capillary forces. These ethanol molecules evaporate after passing an average distance of *L*. Thus a steady capillary flow of ethanol forms as in the case of an ethanol burner. This is why a  $V_{oc}$  can be obtained as long as there is ethanol in the beaker. d) Once the SWNT is pretreated with water, the flowing velocity of ethanol is increased by the Marangoni effect. Both the inducing rate of  $V_{oc}$  and the value of  $V_{oc}$  are enhanced. The  $V_{oc}$  decreases gradually to a value similar to that for pure ethanol after the water molecules evaporate completely.

velocity of ethanol (v) can be calculated using Equation (1):

$$\nu = r_{\rm c}^2 \left( P_{\rm C} - P_L \right) / 8\eta L \tag{1}$$

where  $\eta$  is the viscosity of the liquid,  $r_c$  is the average radius of the capillary channels,  $P_c$  is the capillary pressure difference at the ethanol level in the beaker across the interface  $(2\sigma/r_c)$ , where  $\sigma$  is the surface tension),  $P_L$  is the hydrostatic pressure of liquid ( $P_L = \rho g L \sin \alpha$ , where  $\rho$  the density of liquid, g is the gravitational constant). The velocity of ethanol is estimated to be 57  $\mu$ m·s<sup>-1</sup>. Therefore, the time for the ethanol to rise across L is  $L/\nu = 14.1 \text{ mm}/(57 \,\mu$ m·s<sup>-1</sup>) = 247 s, which is in good consistency with that found in Figure 3a.

If the beaker is covered, the evaporation of ethanol will increase the vapor pressure inside the beaker. When the vapor pressure finally reaches the saturated pressure, the dynamic, directional flow of ethanol will be stopped, and the induced  $V_{\rm oc}$  decreases back to its original value as shown in Figure 2b.

When the length of the SWNT rope is longer than *L*, the added ethanol just raises the level of ethanol in the beaker and that inside the channels among the SWNTs (Fig. 3c). In this situation,

no change in  $V_{\rm oc}$  can be observed (Fig. 2c, time point 1). As a certain amount of ethanol is added and the length of SWNT rope is equal to or less than *L*, the  $V_{\rm oc}$  begins to decrease because there are less moving ethanol molecules inside the channels (Fig. 2c, time points 2–6). Finally, the  $V_{\rm oc}$  decreases to zero and does not change anymore as the ethanol level reaches the top electrode (Fig. 2c, time points 7 and 8).

If the SWNT rope is pretreated with water, as ethanol has a lower surface tension than water, the surface tension difference provides an additional driving force for the ethanol climbing up along the channels among the SWNTs, which is called the Marangoni Effect (Fig. 3d).<sup>[19,28,29]</sup> It is rather complicated to calculate the exact enhanced velocity, and the average velocity of ethanol molecules can be estimated as follows.

At present it is unclear whether L still exists when a Marangoni effect enhances the flow of ethanol. Since the induced  $V_{oc}$  is closely related to the interface area between moving ethanol molecules and SWNTs, it is reasonable to assume that when the  $V_{\rm oc}$  reaches the maximum value, the ethanol molecules moved from the bottom electrode and reached the top electrode. From the dynamic characteristics in the inset of Figure 3b, an average velocity of  $\sim 25.0 \,\mathrm{mm\cdot s^{-1}}$  for ethanol can be obtained. Compared to the velocity of  $57 \,\mu m \, s^{-1}$  for pure ethanol, an enhanced factor of 439 for the velocity has been obtained due to the Marangoni effect.<sup>[10]</sup> The inducing rate for the observed  $V_{\rm oc}$  is enhanced by a factor of 300, which is similar to the

enhancing factor (439) of the velocity of ethanol in the pretreated ropes of SWNTs. Meanwhile, the magnitude of  $V_{oc}$  is enhanced by a factor of 3.9 for the water-pretreated SWNT ropes (the corresponding power is increased by a factor of 15). It is reasonable that the  $V_{oc}$  depends both on the velocity of ethanol and the number of ethanol molecules at the SWNT/water (or ethanol) interface, both of which vary with time. When the  $V_{oc}$  reaches its maximum value (Fig. 3b), the velocity of ethanol decreases and the number of ethanol molecules at the SWNT/water (or ethanol) interface increases. This joint effect makes the  $V_{oc}$  stable at the maximum value for tens of seconds. The  $V_{oc}$  decreases to a similar value to that of pure ethanol when the water effect disappears.

Because the flow of ethanol inside the channels is driven by surface tension, these generators are named surface-energy generators (SEGs). The electrical power of the SEGs can be estimated by

$$P = V_{\rm oc}^2 / R \tag{2}$$

where *P* is the electrical power, and *R* is the inner resistance. Values of 794 pW and 12.0 nW can be obtained for the



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pure-ethanol SEG and the Marangoni-enhanced SEG (angle  $40^{\circ}$ ), respectively, indicating an enhancing factor of 15 for power by the Marangoni effect. The surface area of the SWNT rope contacting the ethanol can be calculated by  $\pi r^2 / \sin \alpha$ , where *r* is the diameter of the SWNT rope (~0.6 mm). The optimum powers obtained are 1.8 mW and 27.3 mW per square meter for the pure-ethanol SEG and the Marangoni-enhanced SEG, respectively, which are high enough to drive nanodevices.<sup>[18]</sup> Here we demonstrate that these SEGs can be used to drive a negative-temperature-coefficient thermistor without additional power (Fig. 4 and Supporting Information, Figs. S4-6). The voltage drop across the thermistor (powered solely by the SEG) as a function of time (Fig. 4) shows a reversible increase (or decrease) as the temperature of the thermistor is decreased (or increased). These processes can be repeated and the voltage drop is consistent with the expected changes in the resistance of the thermistor (Supporting Information, Fig. S6).

Control experiments were carried out with films of SWNTs and carbon fibers. A much smaller  $V_{oc}$  is found in SWNT films, indicating that the alignment is beneficial to the formation of oriented and regular channels among the SWNTs and important for the generation of a higher  $V_{oc}$  (Supporting Information, Figs. S7 and S8) and no  $V_{oc}$  is found for carbon fibers (Supporting Information, Fig. S9). The inducing characteristics of  $V_{oc}$  have also been studied for other organic solvents, such as methanol and acetone. These solvents can also induce a  $V_{oc}$  and the  $V_{oc}$  can be enhanced by pretreating of SWNT rope with water, indicating the universality of these phenomena (Supporting Information, Fig. S10).

In summary, we demonstrate that an effective design of SWNTs can be used to convert the surface energy of liquids into electricity. The mechanism could be ascribed to the unique



**Figure 4.** Self-powered system with a SEG. A negative temperature coefficient thermistor is powered by an SEG of SWNTs ( $V_{oc} = 382 \,\mu$ V,  $I_{sc} = 1.42 \,\mu$ A ( $I_{sc}$  is the short-circuit current),  $R = 269 \,\Omega$  (R is the inner resistance). The voltage drop across the thermistor is detected with varying temperature. The hollow arrows indicate the time when the temperature begins to increase or decrease. The solid arrow shows that the voltage becomes stable when temperature keeps constant at 65 °C. Inset: circuit schematic.



channels among individual SWNTs, in which a continuous, steady flux of liquid forms. The inducing rate of  $V_{oc}$ , the magnitude of  $V_{oc}$ , and the optimum output power can be significantly enhanced by the Marangoni effect. These kinds of SEGs can be used to serve as a self-powered system and have a quite different operating mechanism and show the advantages of a smaller inner resistance, a lack of moving parts, and no application of an obvious external force.

## Experimental

SWNTs were prepared by floating catalytic chemical vapor deposition [30]. After growing for about 6 h, large-sale SWNT films with areas as large as 15 cm  $\times$  2 cm could be carefully peeled off from the inside wall of the guartz tube. Diamond wire drawing dies were used to fabricate aligned SWNT ropes as shown in the Supporting Information, Figure S1. In order to thread the SWNT films through the dies, an as-grown SWNT film was firstly immersed into deionized water. A U-shape copper wire (100-µm diameter) was then used to pull the SWNT film out of the water and thread it through the first die. The pulling processes through the other two dies are similar to the first one. For this work, the SWNT samples are treated with diamond for  $\sim$ 2–3 times and the density of the SWNT rope was much smaller, which was different with our previous work [20]. During the experiments, the sample was suspended on a glass slide with both ends connected to the aluminum electrodes. A Keithley S4200 was employed to monitor the induced voltage and take current-volatage (I-V) characterizations of the sample during the experiments.

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- [1] W. A. de Heer, Rev. Mod. Phys. 1993, 65, 611.
- [2] U. Diebold, Surf. Sci. Rep. 2003, 48, 53.
- [3] D. L. Hu, J. W. M. Bush, Nature 2005, 437, 733.
- [4] D. L. Hu, B. Chan, J. W. M. Bush, Nature 2003, 424, 663.
- [5] D. A. Tomalia, A. M. Naylor, I. William, A. Goddard, Angew. Chem. Int. Ed. 1990, 29, 138.
- [6] F. Schreiber, Prog. Surf. Sci. 2000, 65, 151.
- [7] B. Hammer, J. K. Nørskov, Adv. Catal. 2000, 45, 71.
- [8] M. Grunze, Science 1999, 283, 41.
- [9] B. S. Gallardo, V. K. Gupta, F. D. Eagerton, L. I. Jong, V. S. Craig, R. R. Shah, N. L. Abbott, *Science* 1999, 283, 57.
- [10] S. Daniel, M. K. Chaudhury, J. C. Chen, Science 2001, 291, 633.
- [11] K. Ichimura, S.-K. Oh, M. Nakagawa, Science 2000, 288, 1624.
- [12] M. K. Chaudhury, G. M. Whitesides, Science 1992, 256, 1539.
- [13] M. A. Burns, C. H. Mastrangelo, T. S. Sammarco, F. P. Man, J. R. Webster, B. N. Johnsons, B. Foerster, D. Jones, Y. Fields, A. R. Kaiser, D. T. Burke, *Proc. Natl. Acad. Sci. U. S. A.* **1996**, *93*, 5556.
- [14] Z. L. Wang, J. H. Song, Science 2006, 312, 242.
- [15] A. I. Hochbaum, R. Chen, R. D. Delgado, W. J. Liang, E. C. Garnett, M. Najarian, A. Majumdar, P. D. Yang, *Nature* **2008**, 451, 163.
- [16] Y. Qin, X. D. Wang, Z. Wang, Nature 2009, 457, 340.







[17] J. Liu, P. Fei, J. H. Song, X. D. Wang, C. S. Lao, R. Tummala, Z. L. Wang,

- Nano Lett. **2008**, *8*, 328. [18] B. Z. Tian, X. L. Zheng, T. J. Kempa, Y. Fang, N. F. Yu, G. H. Yu, J. L. Huang, C. M. Lieber, *Nature* **2007**, *449*, 885.
- [19] L. E. Scriven, C. V. Sternling, Nature 1960, 187, 186.
- [20] G. T. Liu, Y. C. Zhao, K. Deng, Z. Liu, W. G. Chu, J. R. Chen, Y. L. Yang, K. H. Zheng, H. B. Huang, W. J. Ma, L. Song, H. F. Yang, C. Z. Gu, G. H. Rao, C. Wang, S. S. Xie, L. F. Sun, *Nano Lett.* **2008**, *8*, 1071.
- [21] G. T. Liu, Y. C. Zhao, K. H. Zheng, Z. Liu, W. J. Ma, Y. Ren, S. S. Xie, L. F. Sun, *Nano Lett.* **2009**, *9*, 239.
- [22] Y. C. Zhao, L. Song, K. Deng, Z. Liu, Z. X. Zhang, Y. L. Yang, C. Wang, H. F. Yang, A. Z. Jin, Q. Luo, C. Z. Gu, S. S. Xie, L. F. Sun, *Adv. Mater.* **2008**, 20, 1772.

[23] B. R. Goldsmith, J. G. Coroneus, V. R. Khalap, A. A. Kane, G. A. Weiss, P. G. Collins, *Science* **2007**, *315*, 77.

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- [24] S. Ghosh, A. K. Sood, N. Kumar, Science 2003, 299, 1042.
- [25] P. Král, M. Shapiro, Phys. Rev. Lett. 2001, 86, 131.
- [26] M. Majumder, N. Chopra, R. Andrews, B. J. Hinds, Nature 2005, 438, 44.
- [27] J. K. Holt, H. G. Park, Y. Wang, M. Stadermann, A. B. Artyukhin, C. P. Grigoropoulos, A. Noy, O. Bakajin, *Science* **2006**, *312*, 1034.
- [28] Y. J. Cai, B.-m. Zhang Newby, J. Am. Chem. Soc. 2008, 130, 6076.
- [29] A. M. Cazabat, F. Heslot, S. M. Troian, P. Carles, Nature 1990, 346, 824.
- [30] W. J. Ma, L. Song, R. Yang, T. H. Zhang, Y. C. Zhao, L. F. Sun, Y. Ren, D. F. Liu, L. F. Liu, J. Shen, Z. X. Zhang, Y. J. Xiang, W. Y. Zhou, S. S. Xie, *Nano Lett.* **2007**, *7*, 2307.

